Engineering Notes

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80 -00 9 Impact of a Rigid Plate on a Half-Strip of Incompressible Liquid

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S UPPOSE the half-strip DOHD of width h (Fig. 1) of an incompressible liquid of density ρ and the rigid walls DO and OH_{∞} are at rest before the moment of time t=0. At the moment t=0 the side wall OH_{∞} is impulsively forced to move in the positive direction on the x-axis and it continues to move during the period of time $0 < t \le \Delta t$. Let the velocity of the wall OH_{∞} at $t=\Delta t$ be u_0 and let the pressure p at the free surface HD be equal to zero.

To obtain the results for velocity distributions v_x and v_y within the liquid at $t = \Delta t$, it is necessary to solve the boundary value problem 1 for the Laplace equation $\nabla^2 \phi = 0$. (ϕ is the velocity potential $v_x = \partial \phi / \partial x$ and $v_y = \partial \phi / \partial y$) with the boundary conditions $v_y(\text{OD}) = \theta$, $v_x(\text{OH}) = u_\theta$, $\phi(\text{HD}) = \theta$, and $\phi(\text{D}) = \theta$.

For the Fourier cosine transformation²

$$\bar{\phi}_c = \sqrt{2/\pi} \int_0^\infty \phi \cdot \cos(\lambda x) \, \mathrm{d}x$$

we have

$$\bar{\phi}_c = -\sqrt{2/\pi} \left(u_0/\lambda^2 \right) \cdot \left\{ I - \left[\cosh\left(\lambda y\right)/\cosh\left(\lambda h\right) \right] \right\}$$

and then

$$(X = \pi x/2h, \quad X^* = \pi x^*/2h, \quad Y = \pi y/2h)$$

$$v_y = \sqrt{2/\pi} \cdot \int_0^\infty (\partial \bar{\phi}_c / \partial y) \cos(\lambda x) \, d\lambda = \frac{u_\theta}{\pi} \ln \frac{\cosh X + \sin Y}{\cosh X - \sin Y}$$

On free surface

$$v_y(\text{HD}) = (2u_0/\pi) \cdot \ln \coth (\pi x/4h)$$

Now

$$v_x = \int_h^y \frac{\partial v_y}{\partial x} \cdot dy = \frac{2u_0}{\pi} \cdot \arctan\left(\frac{\cos Y}{\sinh X}\right)$$

The velocity $v_x(x>0,y=h)=0$, and the point H is the point of discontinuity for v_x !

At the wall OH

$$\phi(0,y) = \int_{+\infty}^{0} v_x dx = -(4u_0 h/\pi^2) D(Y)$$

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Index category: Hydrodynamics.

$$D(Y) = \int_0^{+\infty} \arctan\left(\frac{\cos Y}{\sinh X^*}\right) dX^*$$

$$=2\cos(Y)\cdot\int_0^\infty \frac{X^*\cdot\cosh X^*\cdot\mathrm{d}X^*}{\cosh(2X^*)-\cos(2Y+\pi)}$$

$$= \pi \cdot \ln(2) - 2 \cdot \{L(\pi/4 + Y/2) + L(\pi/4 - Y/2)\}$$

where L is Lobachevskiy's function.³

At the point O the velocity potential $\phi(O) = -(8u_0h/\pi^2)G$, G is Catalan's constant, $^3G = 0.91596...$ In general

$$\phi = -\frac{4u_0h}{\pi^2} \cdot D(Y) \cdot \{1 - \frac{1}{D(Y)} \int_0^{\pi x/2h} \arctan\left(\frac{\cos Y}{\sinh X^*}\right) dX^*\}$$

For the total impulse P of pressure forces on the side wall OH_{∞} for the period of time $0 \le t \le \Delta t$ and for the unit of the length perpendicular to the plane of Fig. 1 we have

$$P(OH_{\infty}) = -\rho \int_{0}^{h} \phi(o,y) dy = \frac{8u_{0}h^{2}\rho}{\pi^{3}} \cdot \int_{Y=0}^{Y=\pi/2} D(Y) dY$$

$$\int_0^{\pi/2} D(Y) dY = (\pi^2/2) \ln(2) - 2(I_1 + I_2)$$

$$I_1 = \int_0^{\pi/2} L\left(\frac{\pi}{4} + \frac{Y}{2}\right) dY; \quad I_2 = \int_0^{\pi/2} L\left(\frac{\pi}{4} - \frac{Y}{2}\right) dY$$

$$I_{I} = -\int_{Y=0}^{Y=\pi/2} dY \int_{\epsilon=0}^{\epsilon=\pi/4} ln \cos\epsilon d\epsilon - \int_{\epsilon=\pi/4}^{\epsilon=\pi/2} ln \cos\epsilon d\epsilon \int_{Y=2\epsilon - (\pi/2)}^{Y=\pi/2} dY$$

$$I_2 = -\int_{\epsilon=0}^{\epsilon=\pi/4} \mathrm{fncosede} \int_{Y=0}^{Y=(\pi/2)-2\epsilon} \mathrm{d} Y$$

$$I_1 + I_2 = \int_0^{\pi/2} (2\epsilon - \pi) \ln \cos \epsilon d\epsilon = -2 \int_0^{\pi/2} \epsilon \ln \sin \epsilon d\epsilon$$

$$P(OH_{\infty}) = \frac{8u_0h^2\rho}{\pi^3} \cdot \left\{ \frac{\pi^2\ln 2}{2} + 4 \cdot \int_0^{\pi/2} \epsilon \ln \sin\epsilon d\epsilon \right\}$$

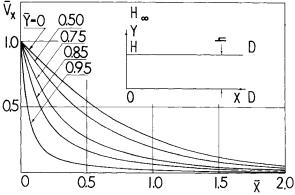


Fig. 1 Dimensionless horizontal velocity.

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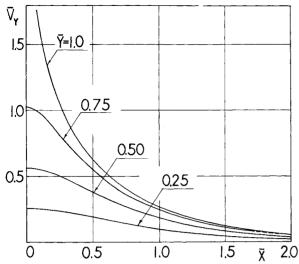


Fig. 2 Dimensionless vertical velocity.

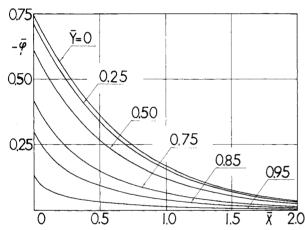


Fig. 3 Dimensionless velocity potential.

In the same manner

$$P(OD) = -\rho \cdot \int_0^\infty \phi(x, \theta) dx$$

$$= \frac{4u_0 h \rho}{\pi^2} \cdot \int_{x=0}^{x=\infty} dx \int_X^\infty \arctan \frac{1}{\sinh X^*} dX^*$$

$$= \frac{4u_0 h \rho}{\pi^2} \int_{X^*=0}^{X^*=\infty} \arctan \left(\frac{1}{\sinh X^*}\right) dX^* \int_0^{x^*} dx = (1/2) \rho u_0 h^2$$

The curves

$$\vec{v}_x = v_x/u_0, \qquad \quad \vec{v}_y = v_y/u_0,$$

and

$$-\bar{\phi} = -\phi/(u_0 h) = (\rho u_0 h)^{-1} \cdot \int_0^{\Delta t} \rho dt$$

as functions of $\bar{x} = x/h$ for some $\bar{y} = y/h$ are given in Figs. 1-3, respectively. Two specific results are: $\bar{\phi}(O) = -0.7425$ and $P(OH)/(\rho u_0 h^2) = 0.5428$.

The influence of impact is mainly observed by the part of the liquid half-strip within the range $0 \le \bar{x} \le 2$.

Acknowledgment

This work was performed while the author was with Payne,

References

¹Lamb, H., *Hydrodynamics*, 6th ed., Dover Publications, New York, 1932.

²Lebedev, N.N., Skalskaya, I.P., and Ufluand Y.S., Worked Problems in Applied Mathematics, Dover Publications, New York, 1979

³Gradshteyn, I.S. and Ryzhick, I.M., Table of Integrals, Series and Products, Academic Press, New York, 1965.

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